

# Imaging of discontinuities in the inverse scattering problem by inversion

verse" GRT. Due to the nature of the asymptotics we are able to give a precise meaning to what is reconstructed by this first-order inversion for arbitrary configurations of sources and receivers, including the case of limited view angles. In particular, we show that the (locations of) *discontin-*

we linearize the relation between the function  $v^{sc}$  and the perturbation  $L_1$ .

Let us now apply this procedure to the Helmholtz equation which describes wave propagation in a fluid of constant density. Without complicating the necessary arguments we

mate solution of (2.6) we use, in place of (2.4), the first term of the geometrical optics approximation

responding transport equations along the rays connecting points  $x, n$  and  $x, \bar{\xi}$ , respectively. In the following sections we



$$\partial X_{\eta}^0(y) = \partial X \setminus \partial X_{\eta}(y).$$

To describe  $\partial X_{\eta}(y)$  and, therefore,  $\partial X_{\eta}^0(y)$  for a given interior point  $y \in X$ , three cases can be distinguished.

(1)  $\hat{n}(y) > \tilde{n}(y)$ . In this case  $\partial X_{\eta}(y)$  is empty since  $\epsilon$  can be chosen sufficiently small

we find that

$$dp = k^{n-1} h(y, \xi) d\xi dk, \tag{4.16}$$

and, thereby,

$$1 = \int_{\partial X_{\eta}(y)} h(y, \xi) d\xi dk \tag{4.17}$$

$$C_\phi = \{(k, \xi, x, y) \in R_+ \times \partial X_\eta^0(y) \times X \times X : \Phi(x, y, \xi) = 0, \\ \nabla_\xi \Phi(x, y, \xi) = 0\}. \quad (4.21)$$

$F$  of the form in (4.27) consists of increasingly smooth pseudodifferential operators. Comparing the first term in the expansion (4.28) with (4.17) we obtain the expansion in (4.19).

Again, consider the Fourier integral operator  $F$  where  $v(k, \xi, \eta)$  is described in (3.1). In many cases of practical

where the bar denotes the complex conjugate. In particular, this relation shows that the Fourier space is covered twice if  $\partial X^0 = \partial X$ . Given the source-receiver configuration of a particular experiment, we can determine the domain in the Fourier space where the function  $\hat{f}$  is known. This domain controls the spatial resolution of the reconstruction. In examples where the domain  $X$  is a half-space (given later in this section) we only have partial coverage since observation points are restricted to the boundary of the half-space. How-

where  $\bar{l}$  and  $\bar{l}_0$  are unit vectors pointing in the direction of  $x_n$  axis and in the direction of the line connecting points  $x$  and  $\eta$ . The generalized backprojection operator  $R^*$  thus is

$$(R^*w)(x) = C_n \int_{\partial X} w(2|x - \eta|, \eta) \bar{l} \cdot \bar{l}_0 d\eta, \quad (6.9)$$

where  $w$  is given in (6.2) or (6.3) depending on the dimension  $n$ . Here we integrate over all source-receiver positions on the boundary  $\partial X$ . The GBO in (6.9) can be considered as a migration scheme with a special weighting function. This migration



$$\phi_{x_1} = \frac{x_1 - \eta - d}{\tilde{r}} + \frac{x_1 - \eta + d}{\tilde{r}},$$

(with the preservation of orientation) to the unit sphere  $S_x^{n-1}$  centered at the origin of the tangent space at the point  $x \in Y$

